



# Auroral processes in satellite data



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# Overview

- Setting the scene
- Auroral processes in the upward current region
- Auroral processes in the downward current region
- Dayside and cusp aurora, theta aurora
- Temporal evolution in the auroral zone

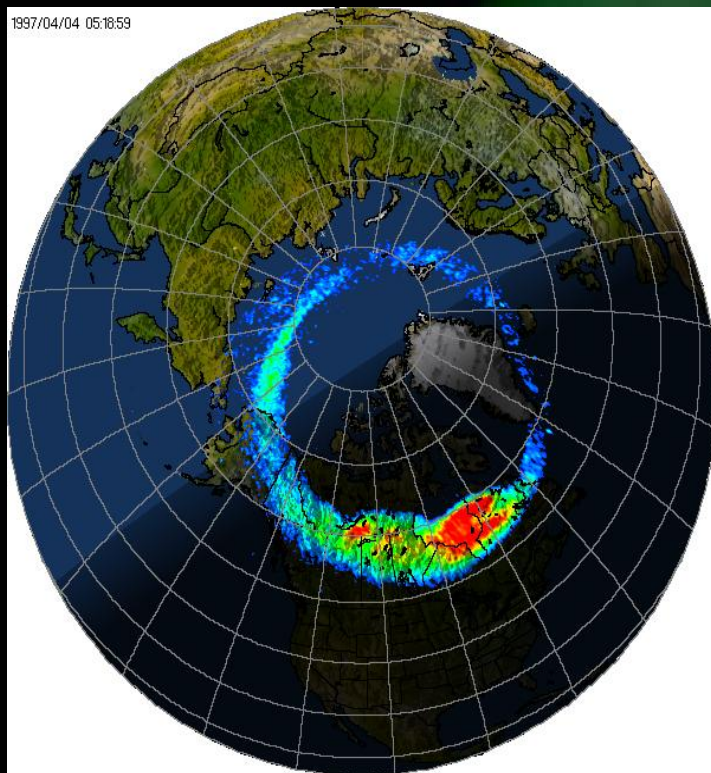
*Very useful reference:*

*Space Science Series of ISSI*  
**Auroral Plasma Physics**

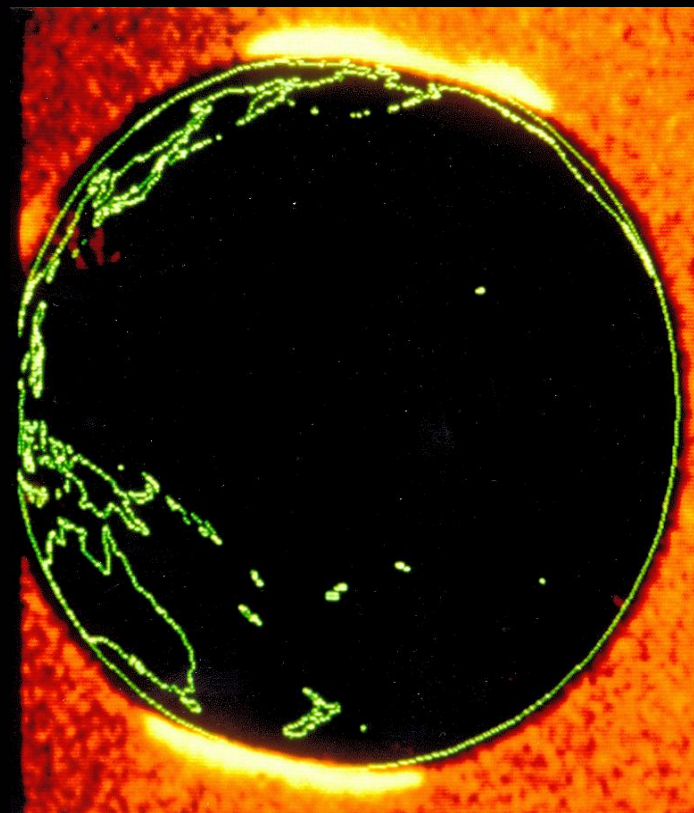
*Götz Paschmann, Stein Haaland, Rudolf Treumann*

(Space Science Reviews, vol 103, 1-4, 2002)

# Auroral ovals

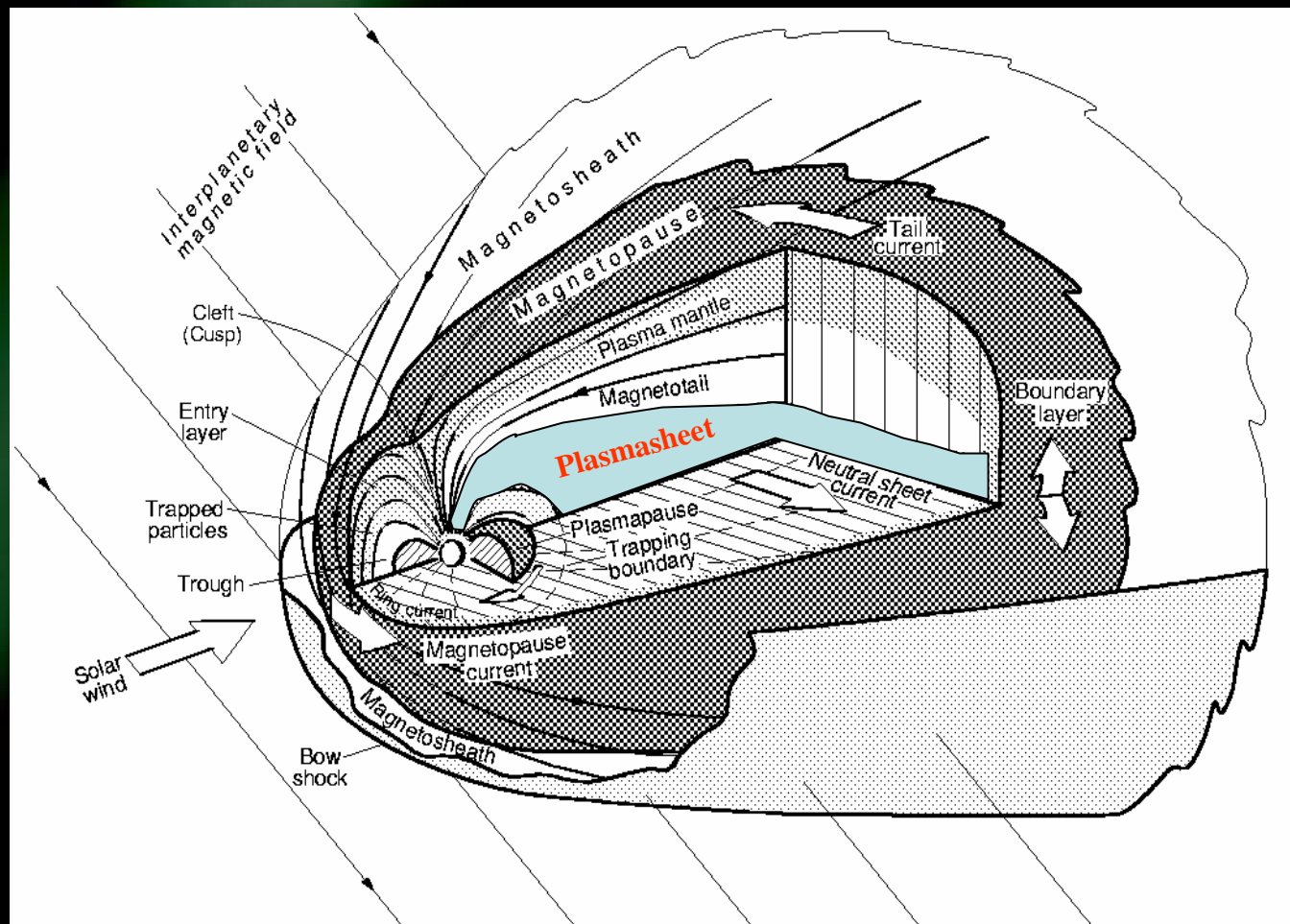


Polar



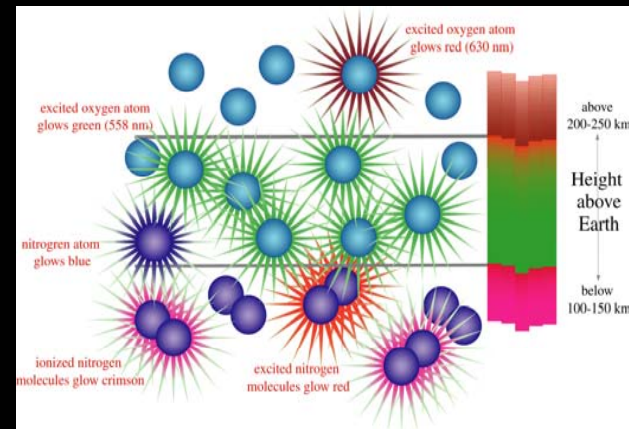
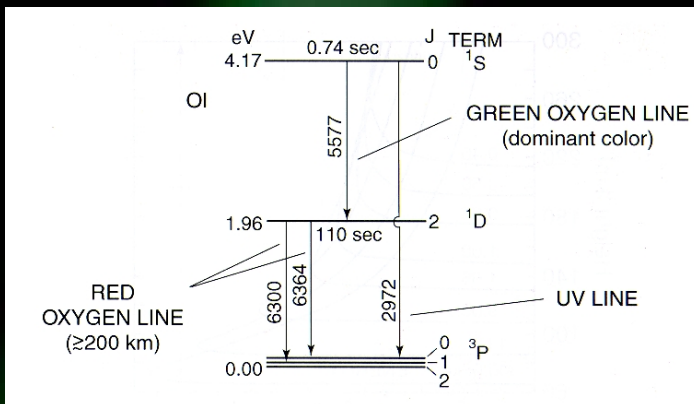
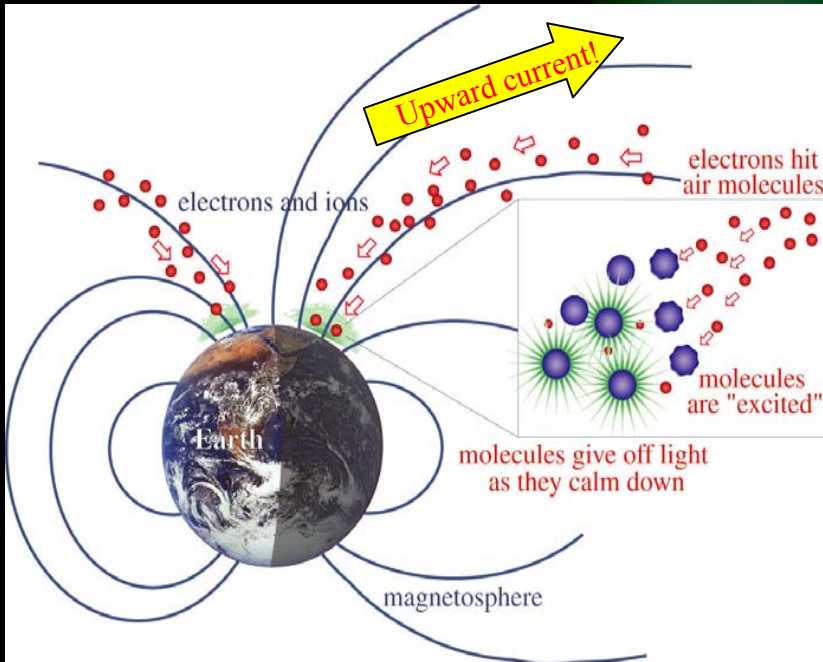
Dynamics Explorer

# The auroral oval is the projection of the plasmasheet onto the atmosphere

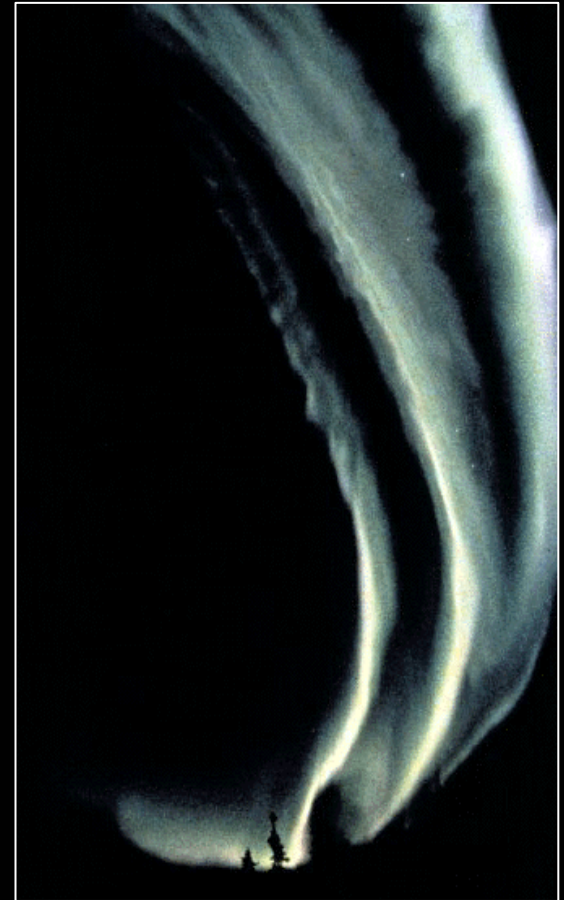




# Auroral emissions

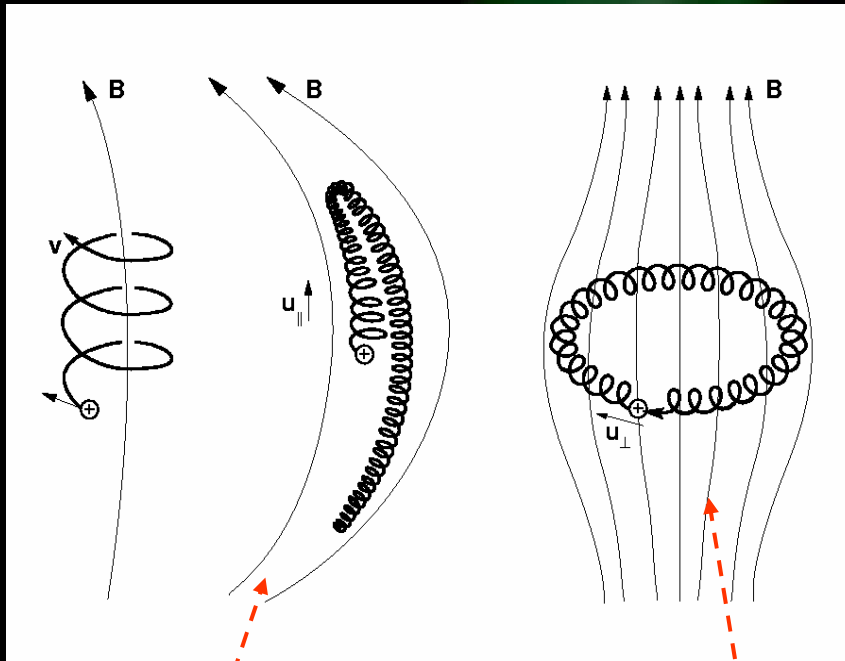


# Homogenous auroral arcs



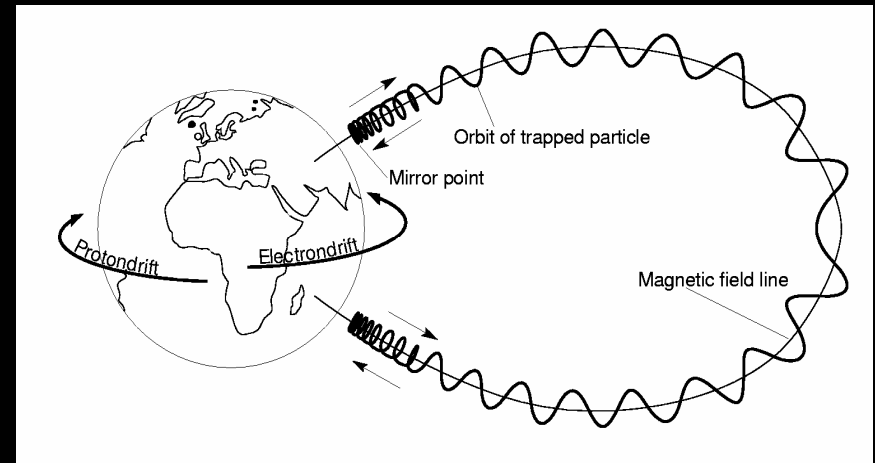
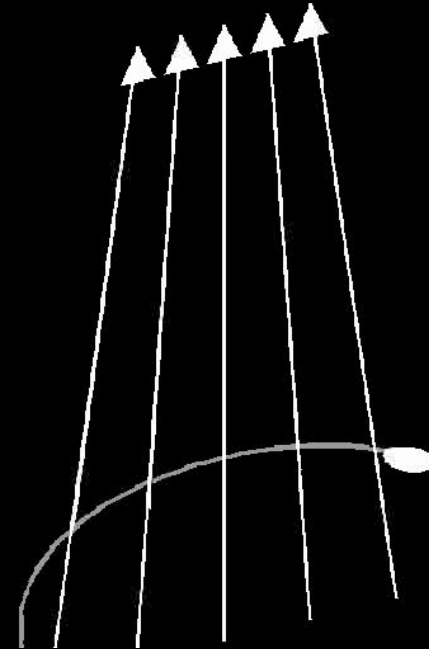
# Particle motion in the geomagnetic field

gyration      longitudinal oscillation      azimuthal drift



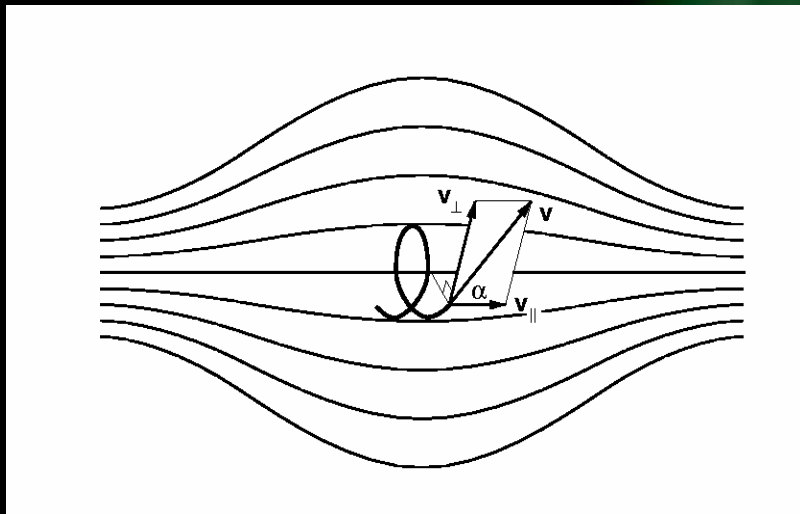
Magnetic mirror

grad B drift





# Magnetic mirror



The magnetic moment  $\mu$  is an *adiabatic invariant*.

$$\mu = \frac{mv_{\perp}^2}{2B} = \frac{mv^2 \sin^2 \alpha}{2B}$$

$mv^2/2$  constant ➡

$$\frac{\sin^2 \alpha}{B} = \text{const}$$

particle turns when  $\alpha = 90^\circ$  ➡

$$B_{\text{turn}} = B / \sin^2 \alpha$$

If maximal B-field is  $B_{\text{max}}$  a particle with pitch angle  $\alpha$  can only be turned around if

$$B_{\text{turn}} = B / \sin^2 \alpha \leq B_{\text{max}} \quad \text{➡}$$

$$\alpha > \alpha_{lc} = \arcsin \sqrt{B / B_{\text{max}}}$$

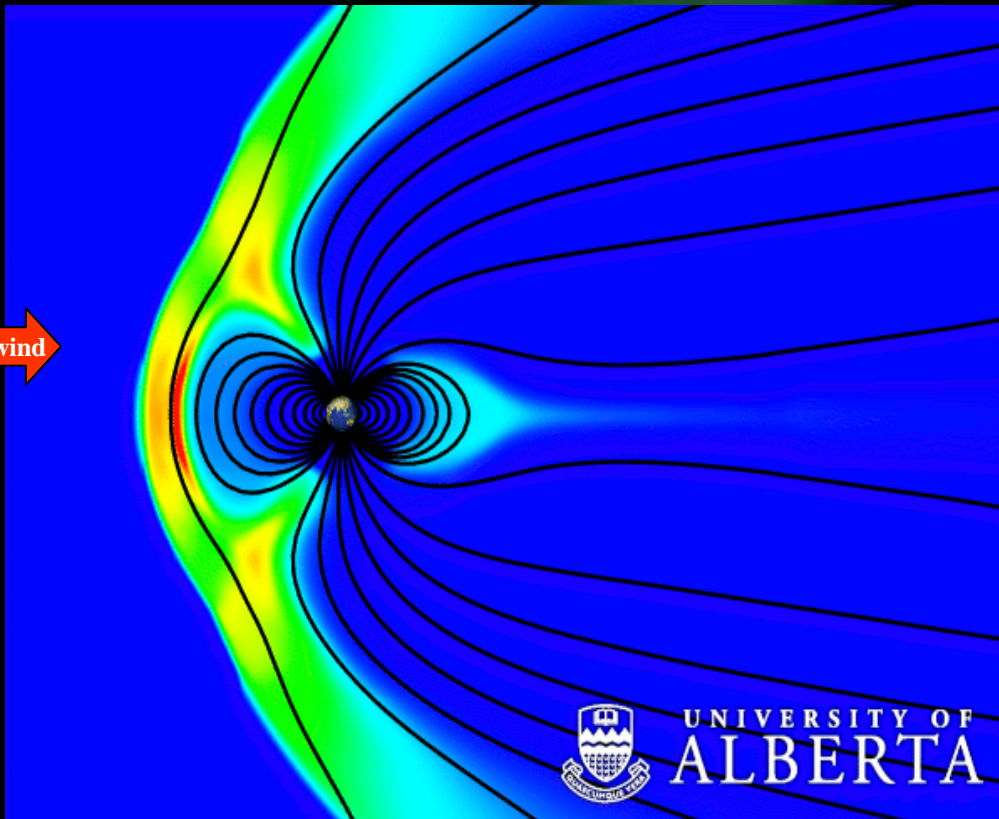
Particles in  
*loss cone* :

$$\alpha < \alpha_{lc}$$

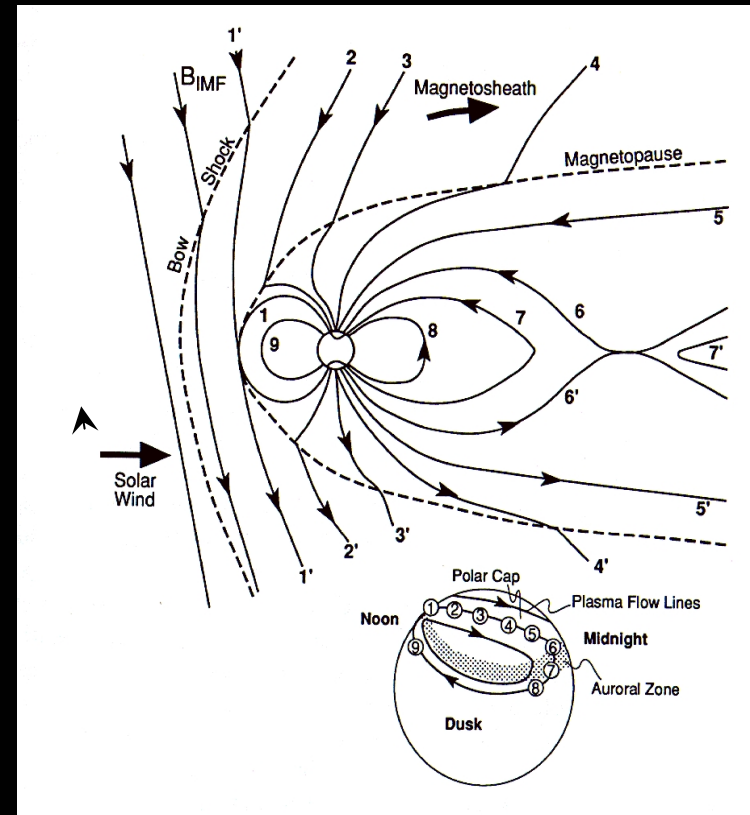
# Magnetospheric convection



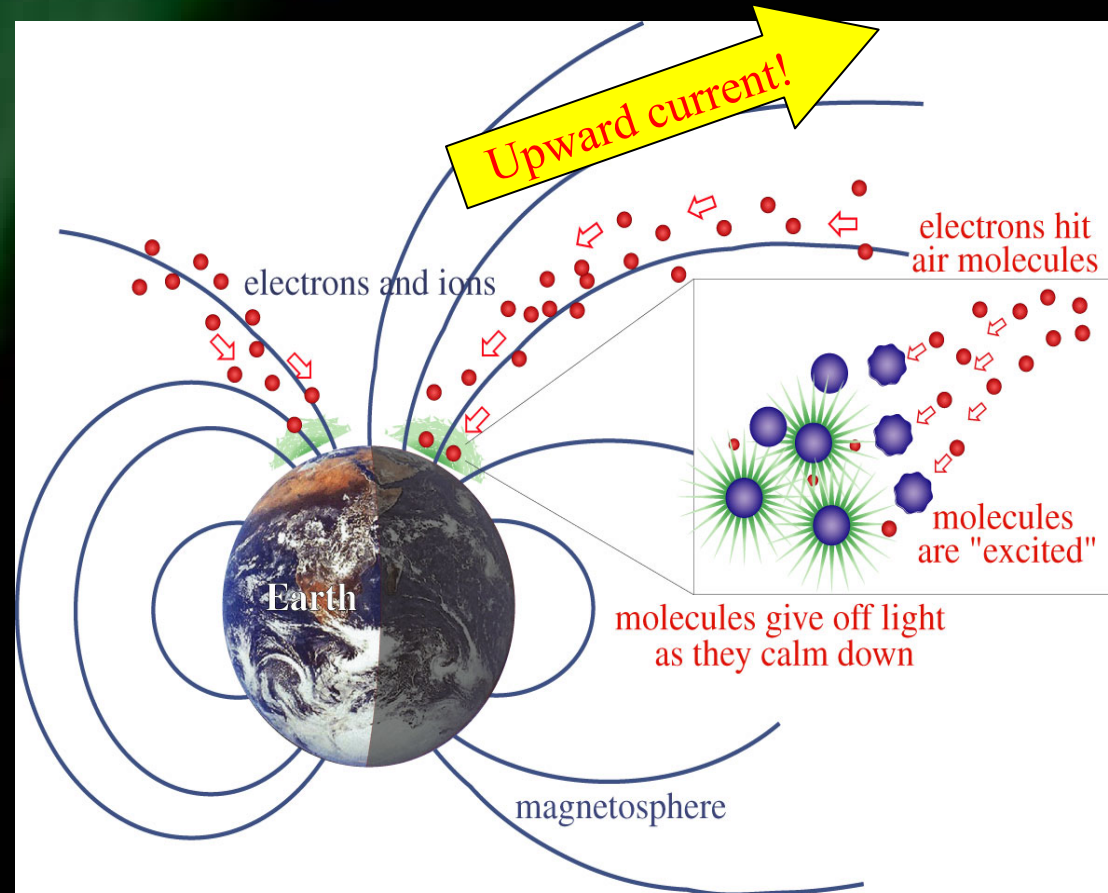
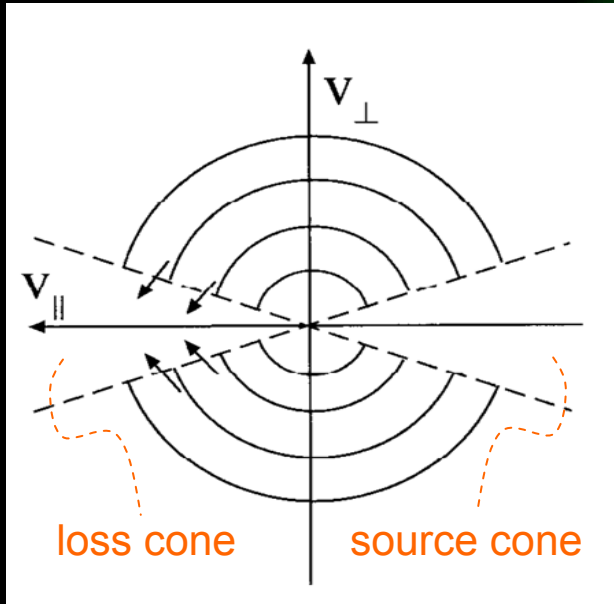
Solar wind



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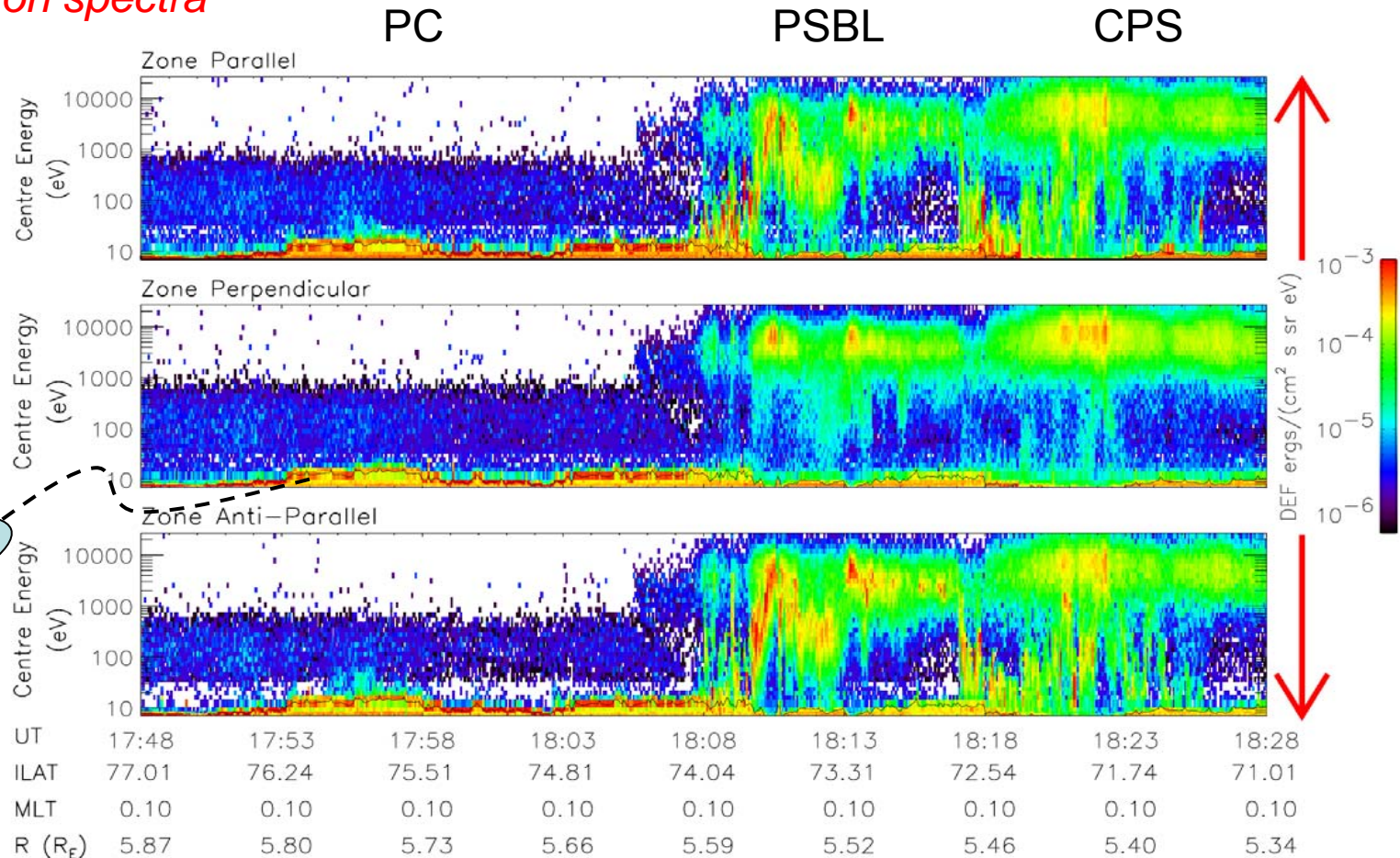


# Atmospheric collisions - emissions



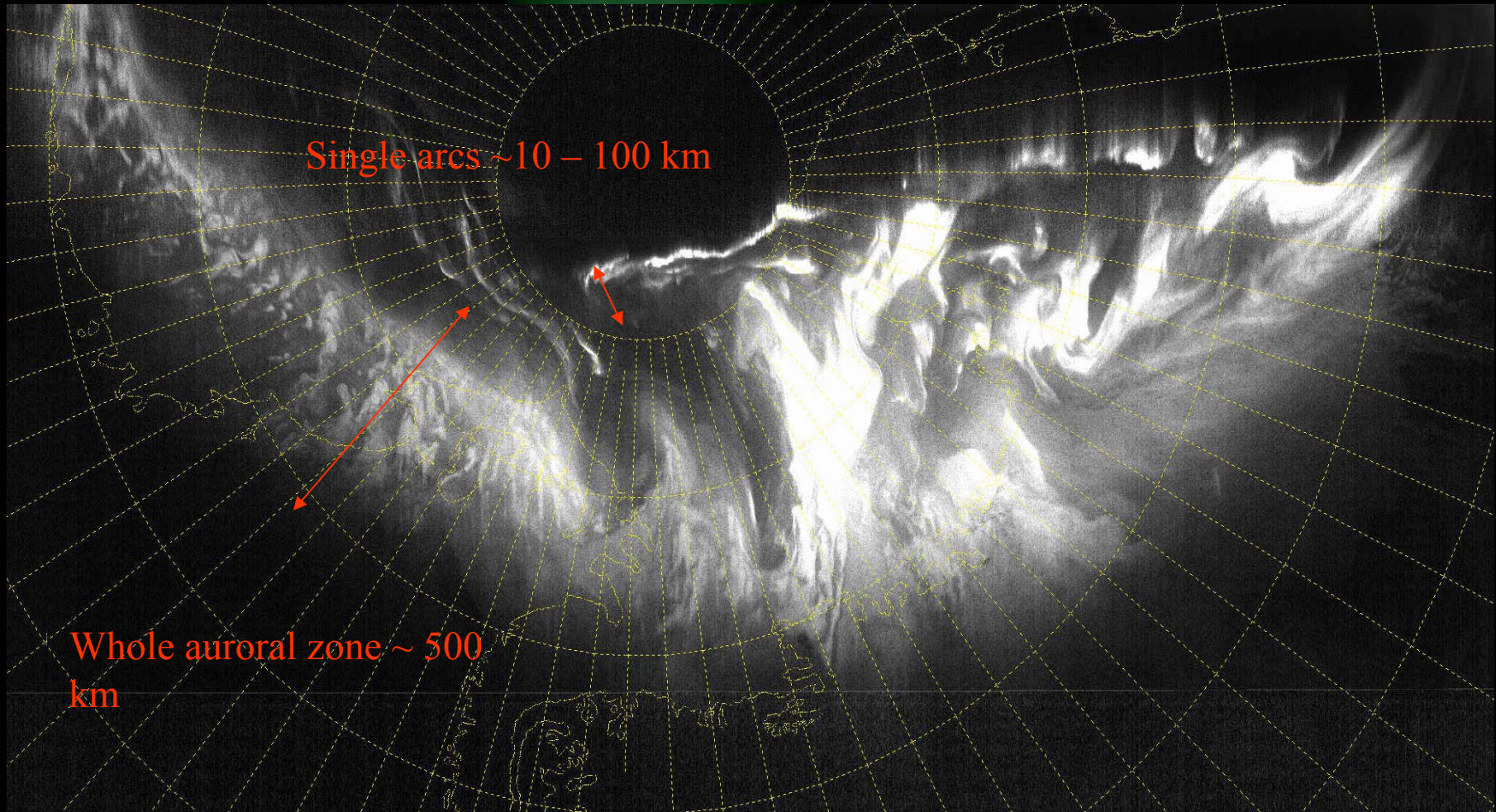
# A typical auroral pass - CLUSTER

## Electron spectra





# Auroral scales



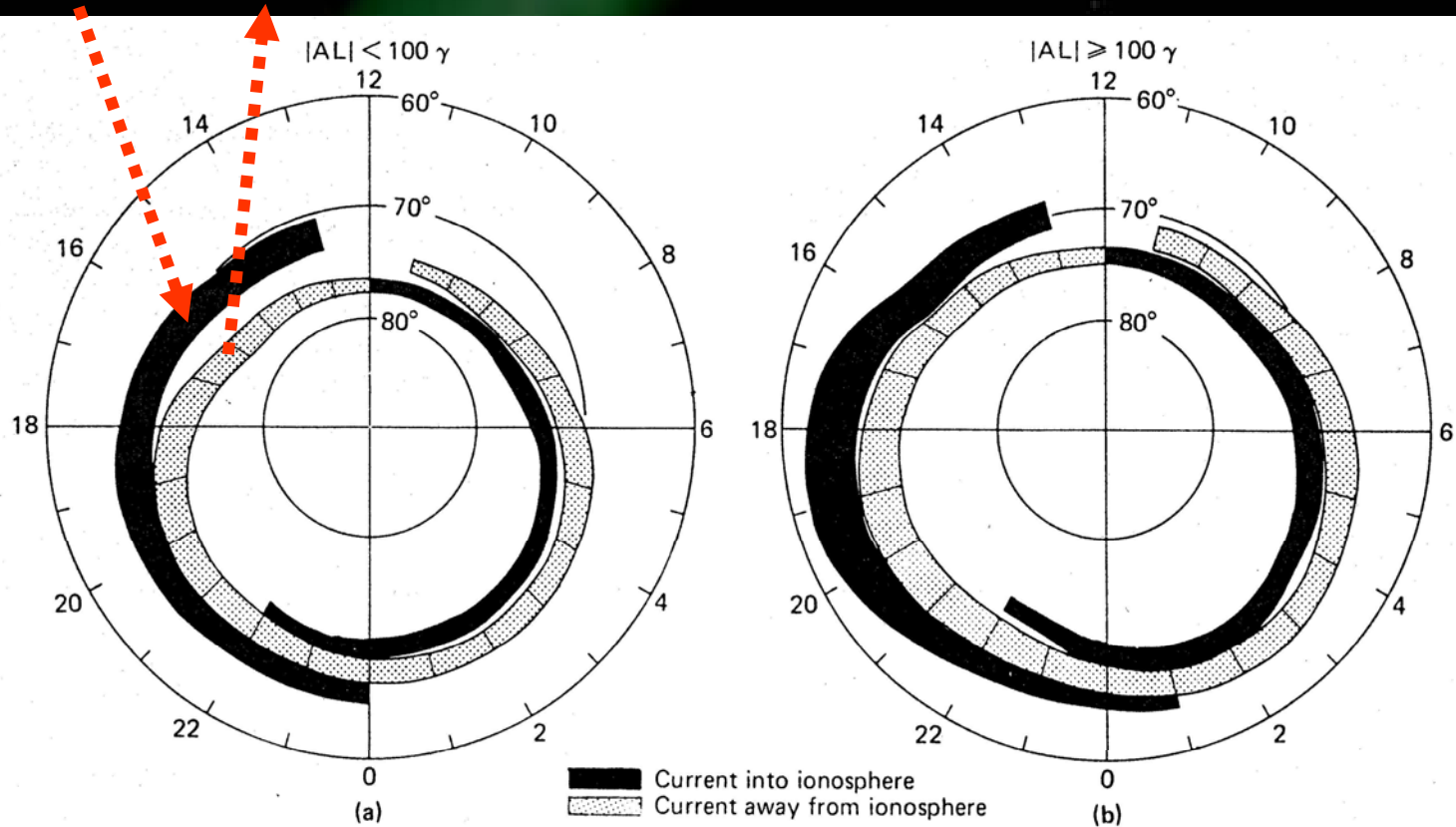
*Photo from DMSP satellite*



# Birkeland currents in the auroral oval

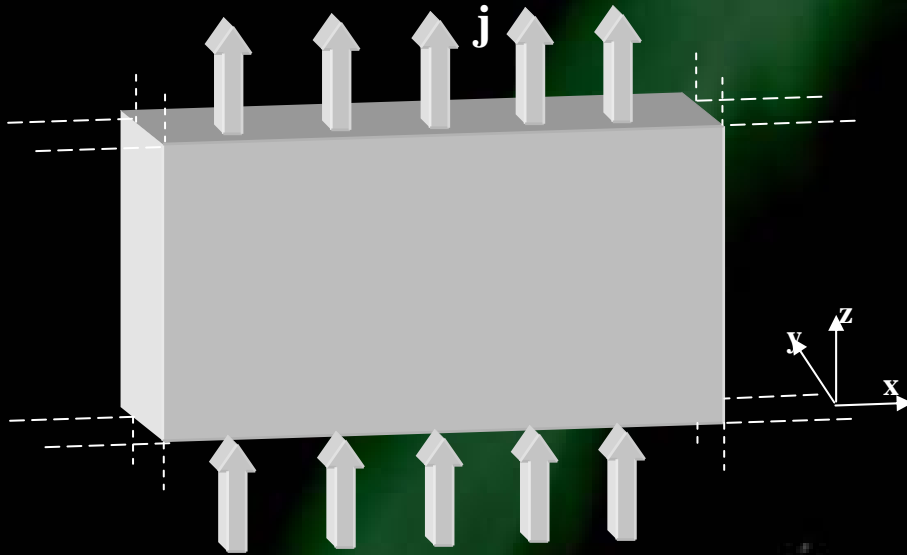
Low geomagnetic activity

High geomagnetic activity





# Current sheet approximation and Ampère's law



$$\left( \frac{\partial B_z}{\partial y} - \frac{\partial B_y}{\partial z}, \frac{\partial B_x}{\partial z} - \frac{\partial B_z}{\partial x}, \frac{\partial B_y}{\partial x} - \frac{\partial B_x}{\partial y} \right) = \mu_0 (j_x, j_y, j_z)$$

But  $\frac{\partial}{\partial x} = 0$  and  $\frac{\partial}{\partial z} = 0$

$$\left( \frac{\partial B_z}{\partial y}, 0, -\frac{\partial B_x}{\partial y} \right) = \mu_0 (0, 0, j_z)$$

Ampère's law (no time dependence):

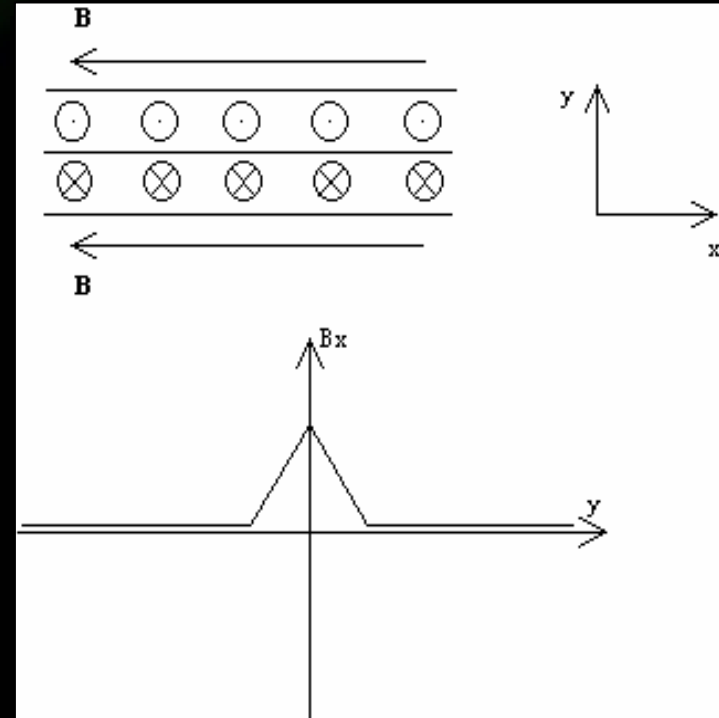
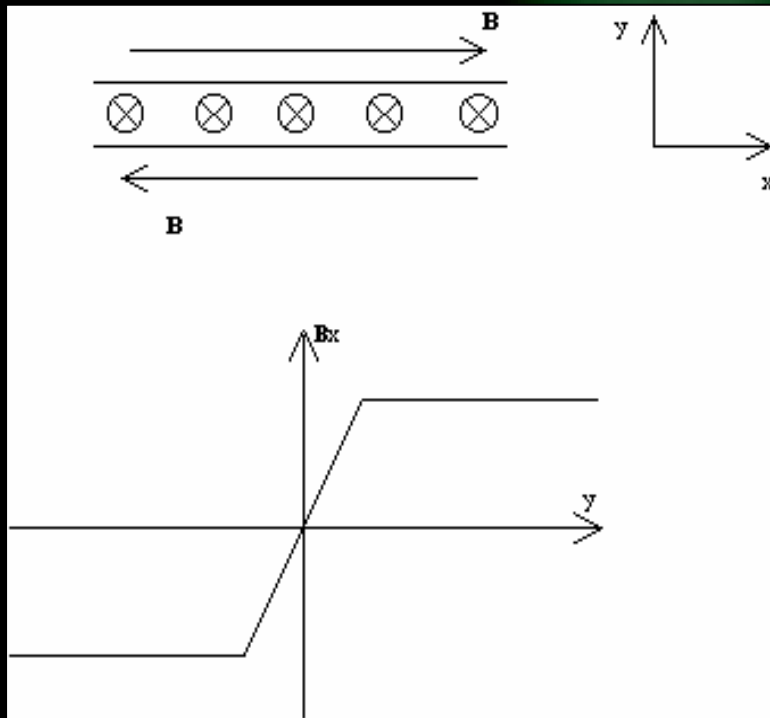
$$\nabla \times \mathbf{B} = \mu_0 \mathbf{j}$$



$$j_z = -\frac{1}{\mu_0} \frac{\partial B_x}{\partial y}$$

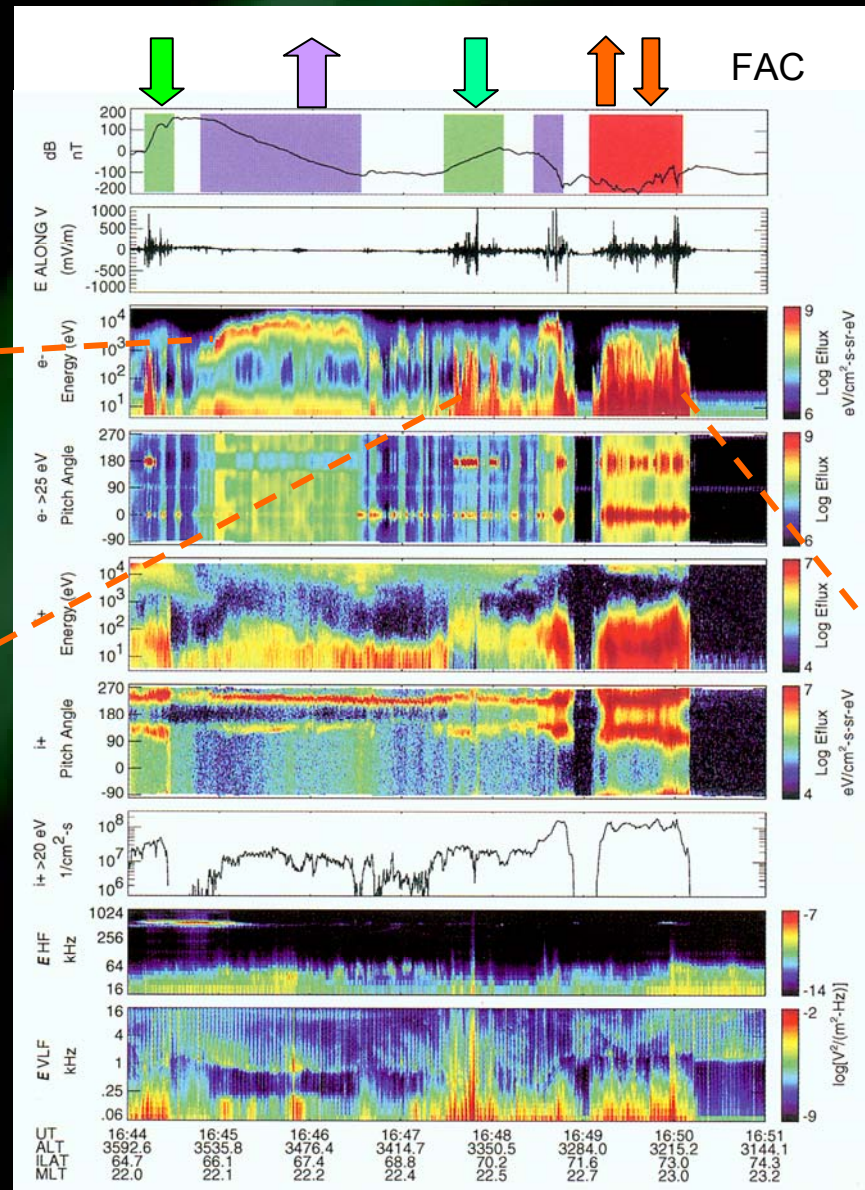
# Current sheet

## Determination of current density by magnetic field measurement



$$j_z = -\frac{1}{\mu_0} \frac{\partial B_x}{\partial y}$$

# Upward and downward current regions



Inverted V

Upgoing electron beams

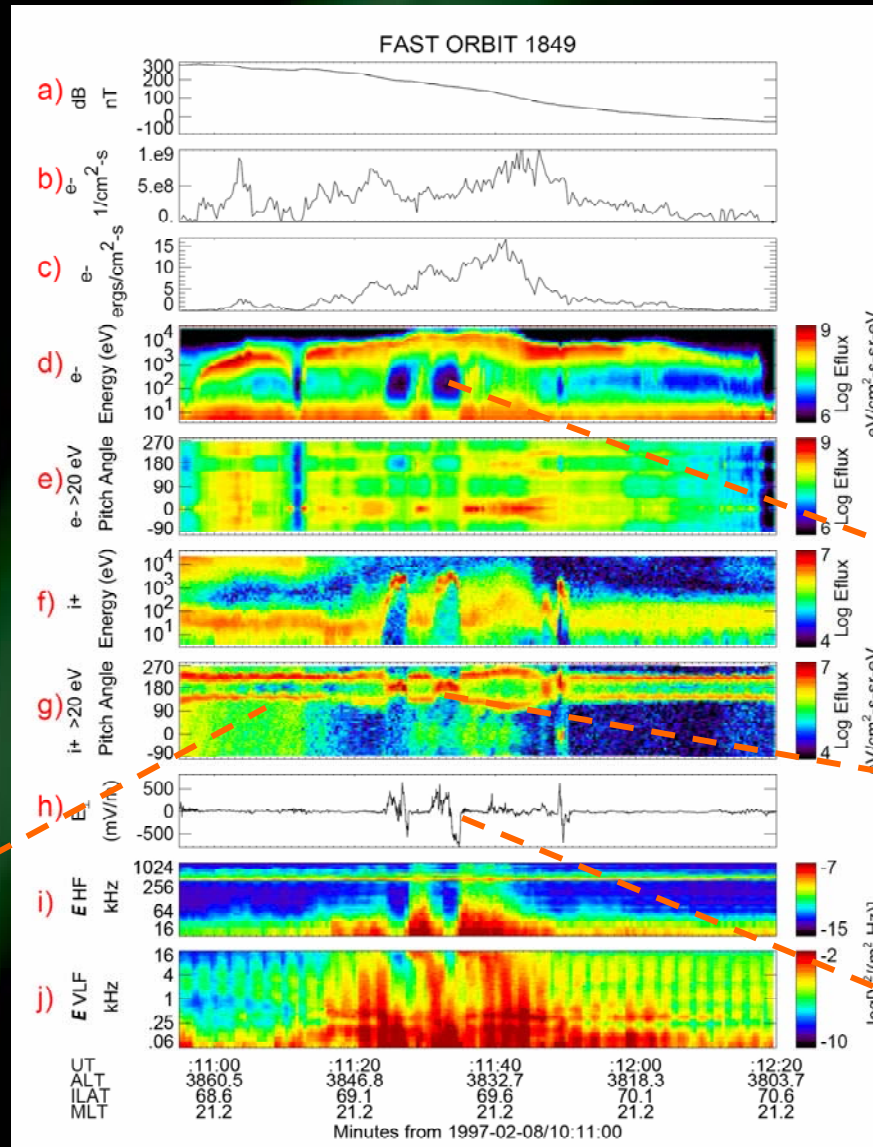
180° = ↑  
0° = ↓

Alfvénic aurora



# Upward current region

## Inverted V arc



Visual limit  $\sim 10^{-3} \text{ Jm}^{-2}\text{s}^{-1}$   
 $= 1 \text{ erg cm}^{-2}\text{s}^{-1}$

$180^\circ = \uparrow$

$0^\circ = \downarrow$

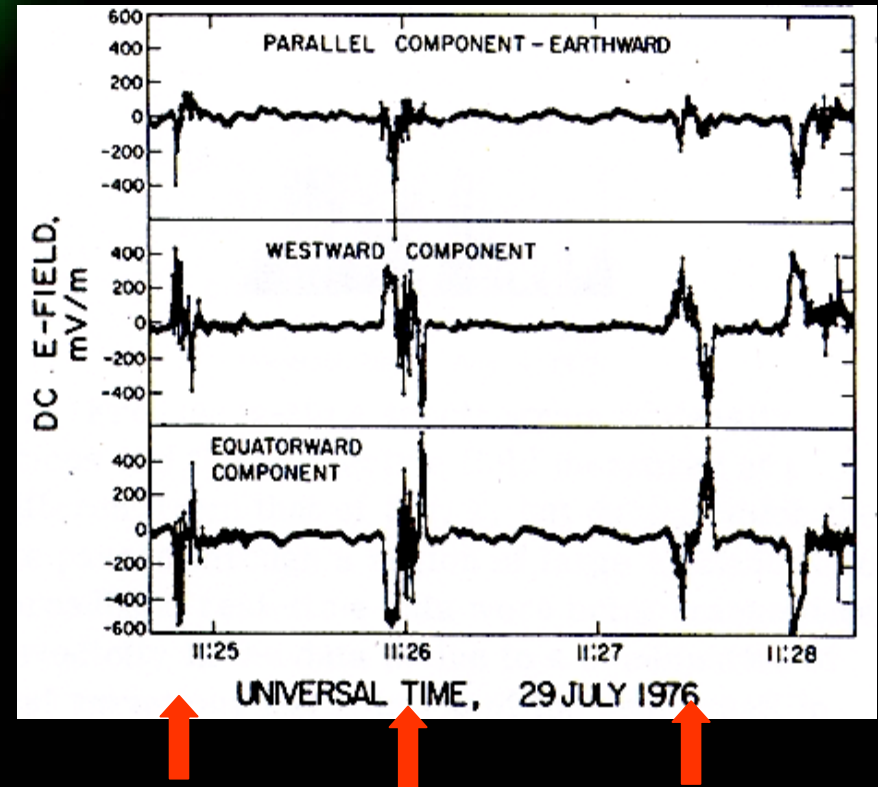
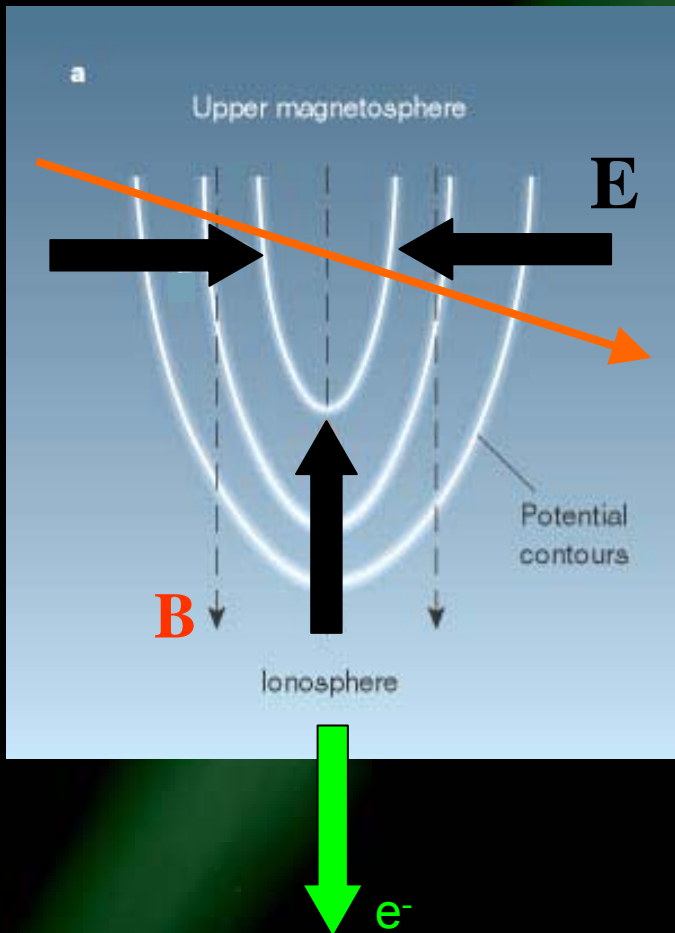
Ion conic

Auroral density cavity

Ion beam

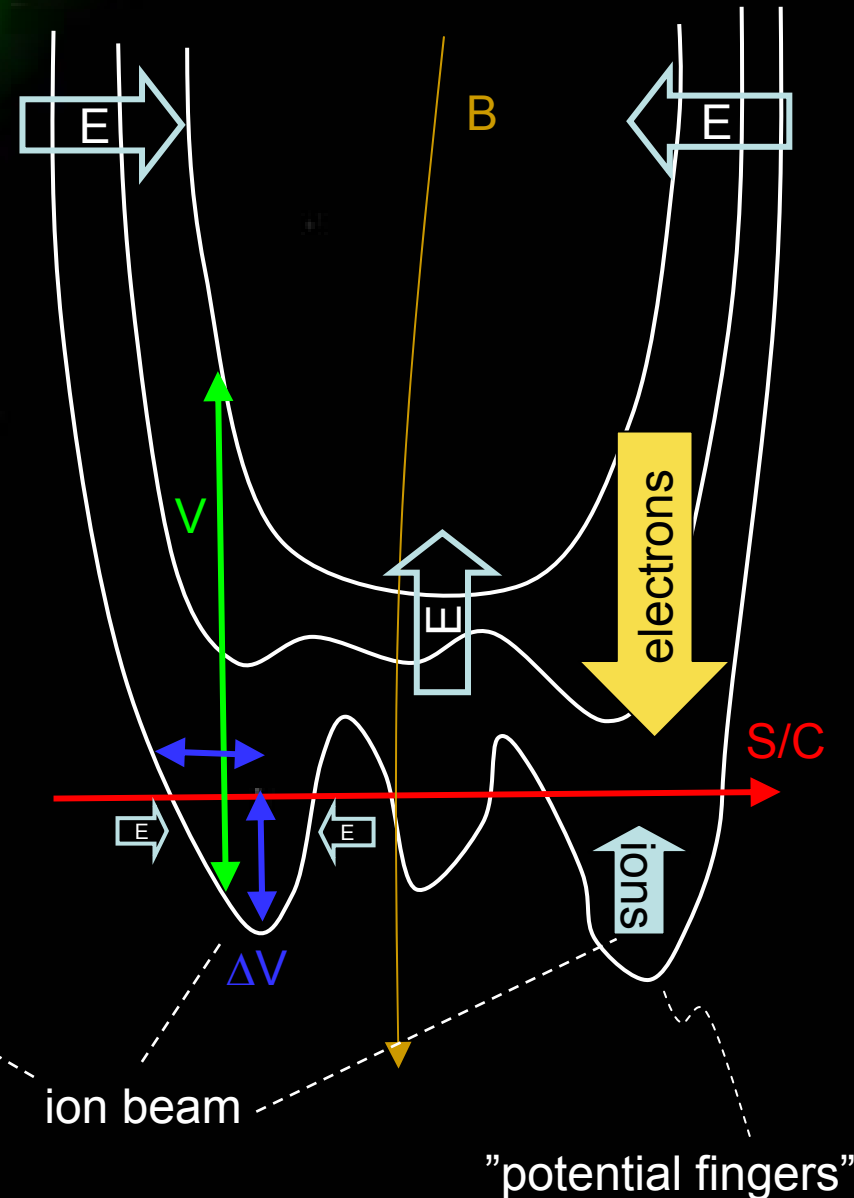
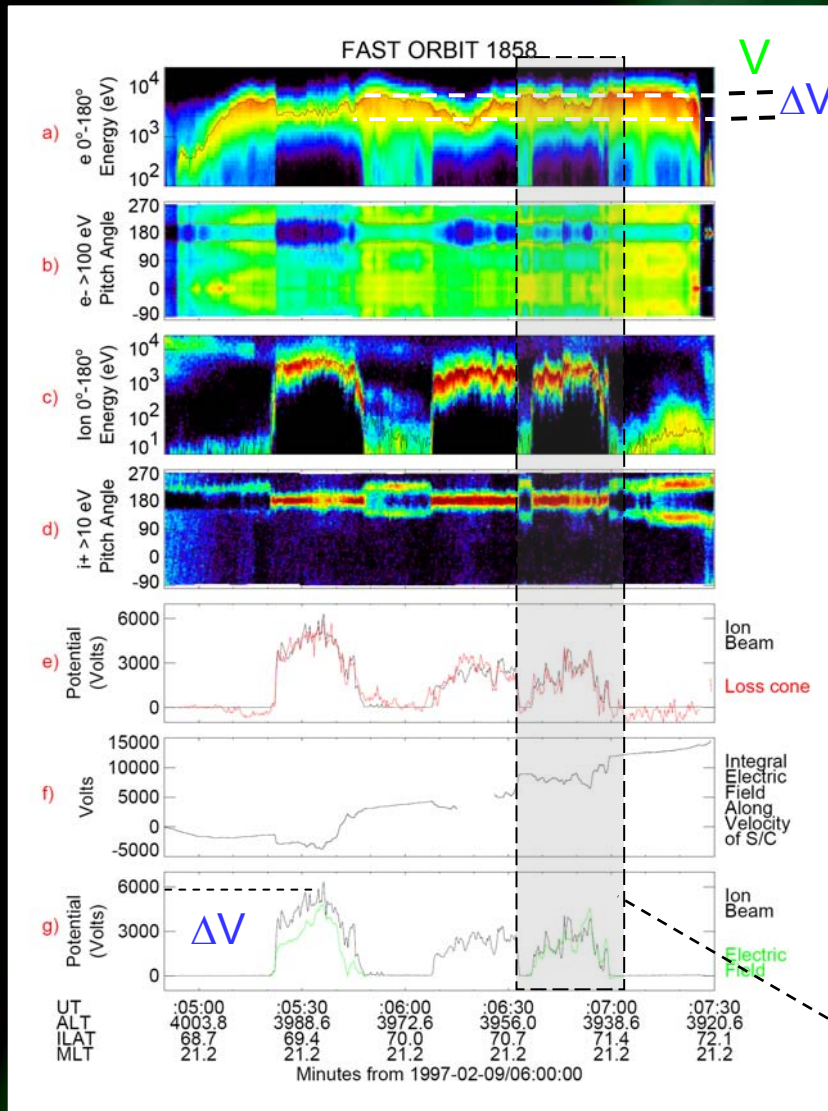
Bipolar electric field structures

# Satellite signatures of U potential

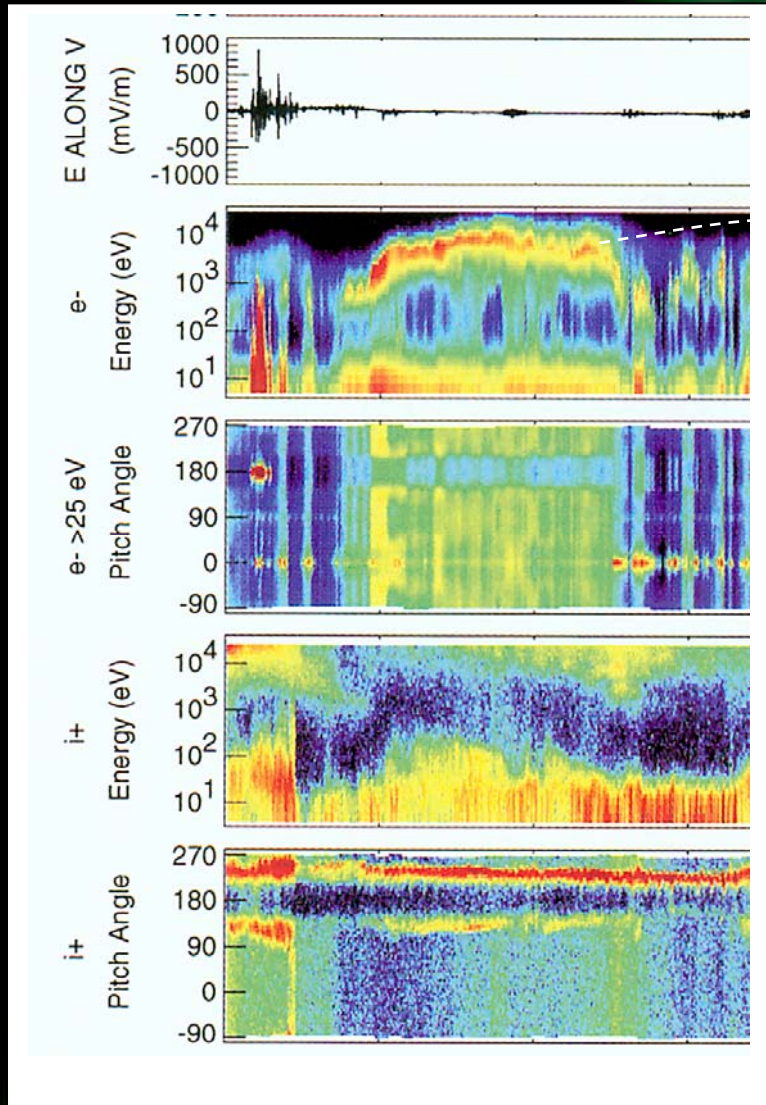


Measurements made by the ISEE satellite (Mozier et al., 1977)

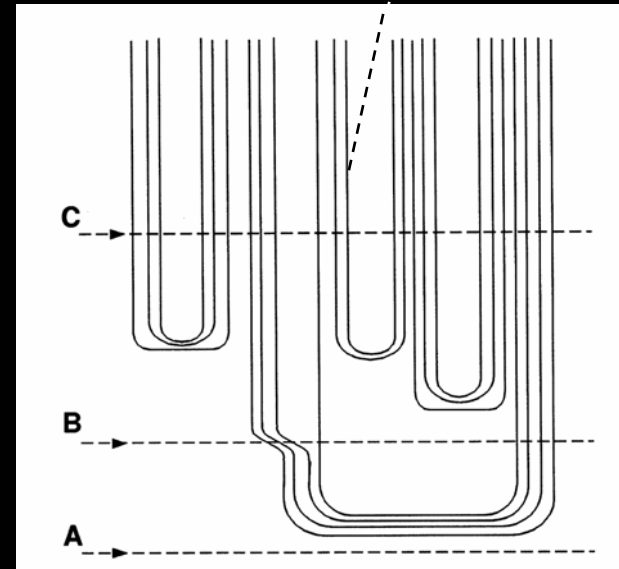
# Acceleration potential structure I



# Acceleration potential structure II



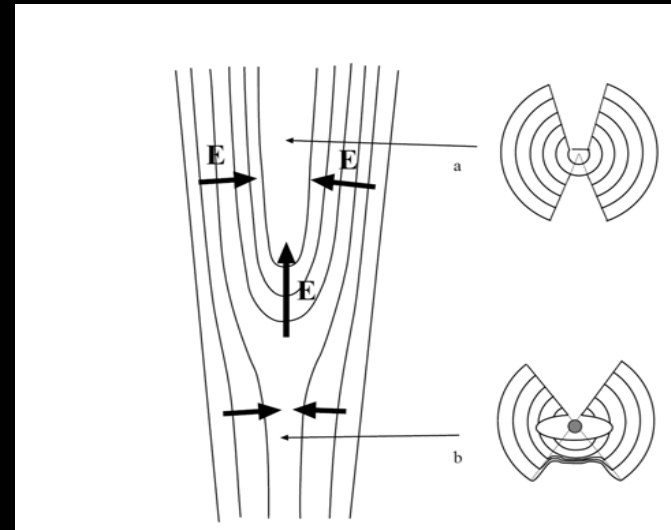
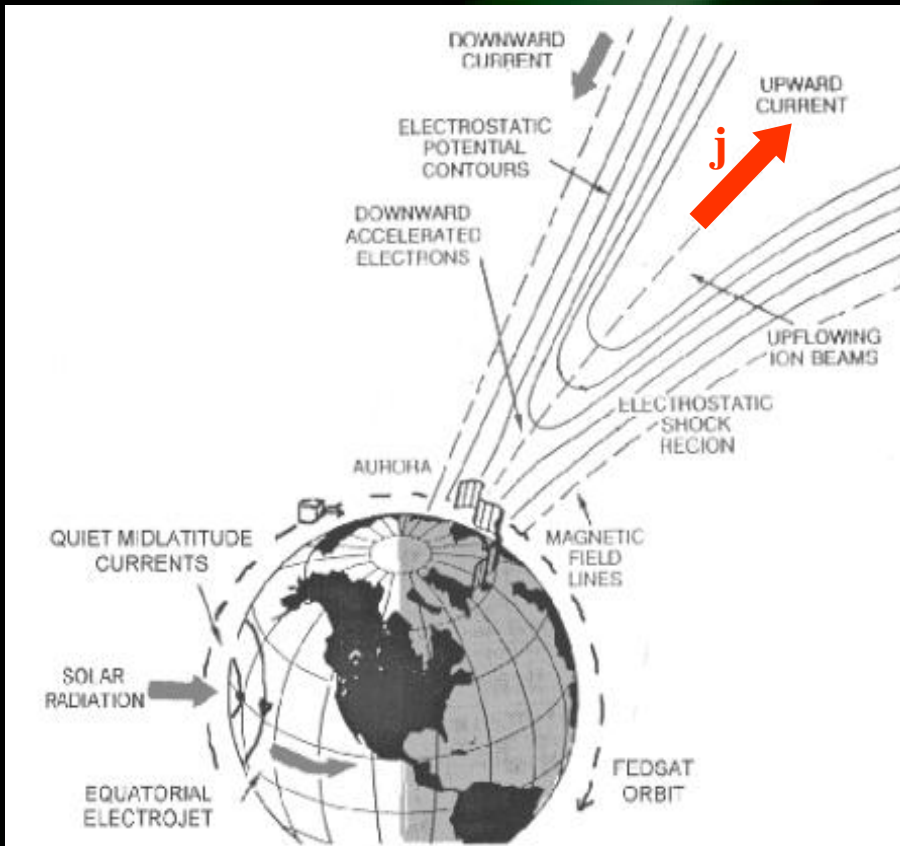
Localized regions of higher energy electrons without associated ion beams





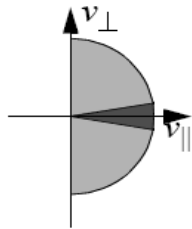
# Why particle acceleration?

- The magnetosphere often acts as a current generator
- Electrons are accelerated downwards by upward E-field.
- This increases the pitch-angle of the electrons, and more electrons can reach the ionosphere, where the current can be closed.



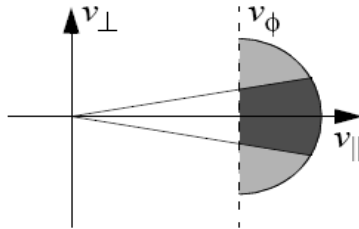


# Auroral currents – Knight relation



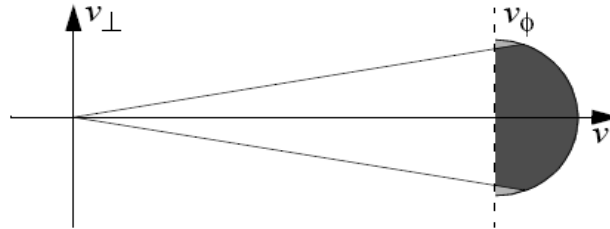
Thermal flow

$$e\Phi_{\parallel} \ll 1$$



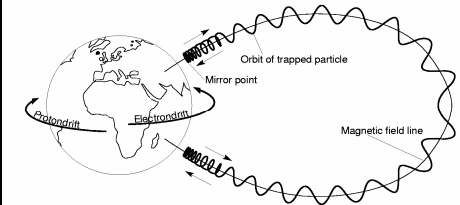
Linear regime

$$1 \ll e\Phi_{\parallel} \ll R_M$$



Saturation

$$R_M \ll e\Phi_{\parallel}$$



Fraction of particles  
in the loss cone:

$$f = \frac{\pi\theta_{lc}}{2\pi} \sim \frac{B_{ms}}{B_{ion}} \sim 2 \times 10^{-4}$$

Thermal current:

$$\begin{aligned} j_{\parallel,ms} &= n_0 e v_{th} f \\ j_{\parallel,ion} &= n_0 e v_{th} f \frac{B_{ms}}{B_{ion}} = \\ &= n_0 e v_{th} \frac{B_{ms}}{B_{ion}} \frac{B_{ion}}{B_{ms}} = n_0 e v_{th} \approx \end{aligned}$$

$$\begin{aligned} [n_e = 0.1 \text{ cm}^{-3}, T_e = 1 \text{ keV}] &\approx \\ &\sim 1 \mu\text{A/m}^2 \end{aligned}$$

Apply a parallel potential drop:

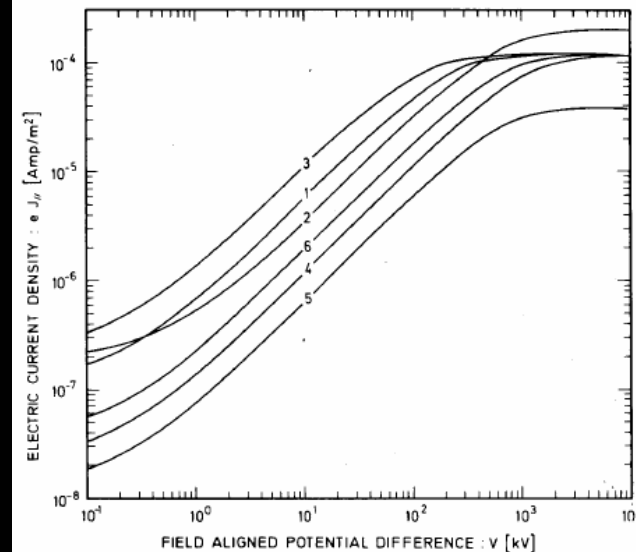
$$j_{\parallel,ion} = n_0 e v_{th} \frac{B_{ion}}{B_{ms}} \left[ 1 - \frac{e^{-xe\Phi_{\parallel}/T_e}}{1+x} \right]$$

Linear regime :

$$j_{\parallel,ion} \approx n_0 e v_{th} \frac{e\Phi_{\parallel}}{k_B T_e} = K \Phi_{\parallel}$$

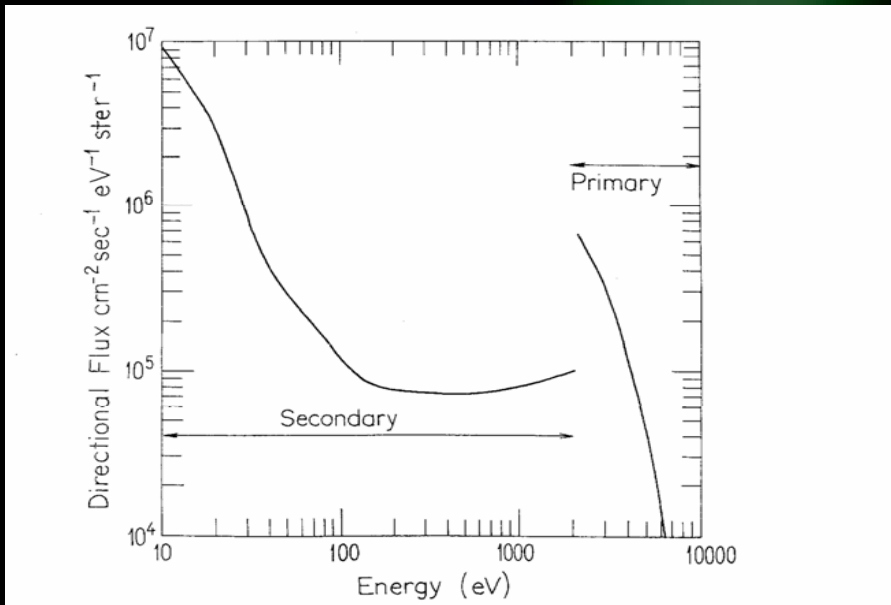
$$K = \frac{e^2 n_0}{\sqrt{2\pi m_e k_B T_e}} \sim 10^{-9} \text{ S/m}^2$$

FRIDMAN AND LEMAIRE: CALCULATION OF AURAL ELECTRON FLUXES

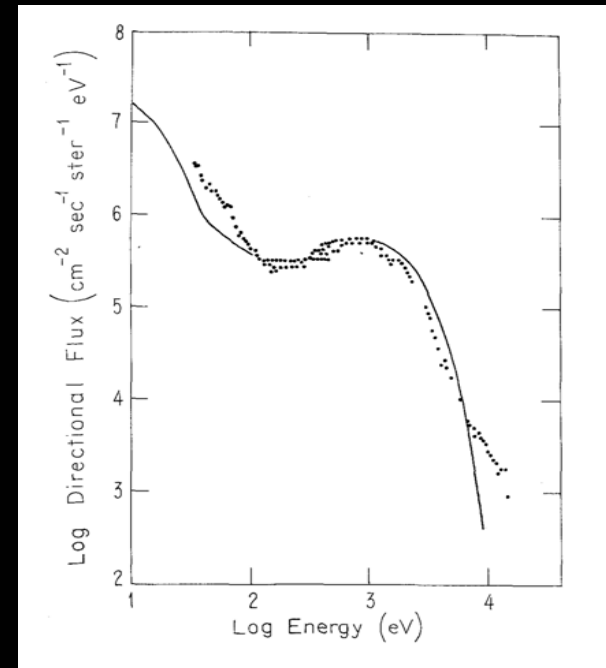


$$v_{th} = \sqrt{T_e / 2\pi m_e} \quad x = \frac{1}{B_l / B_0 - 1}$$

# Particle distributions associated with inverted V's

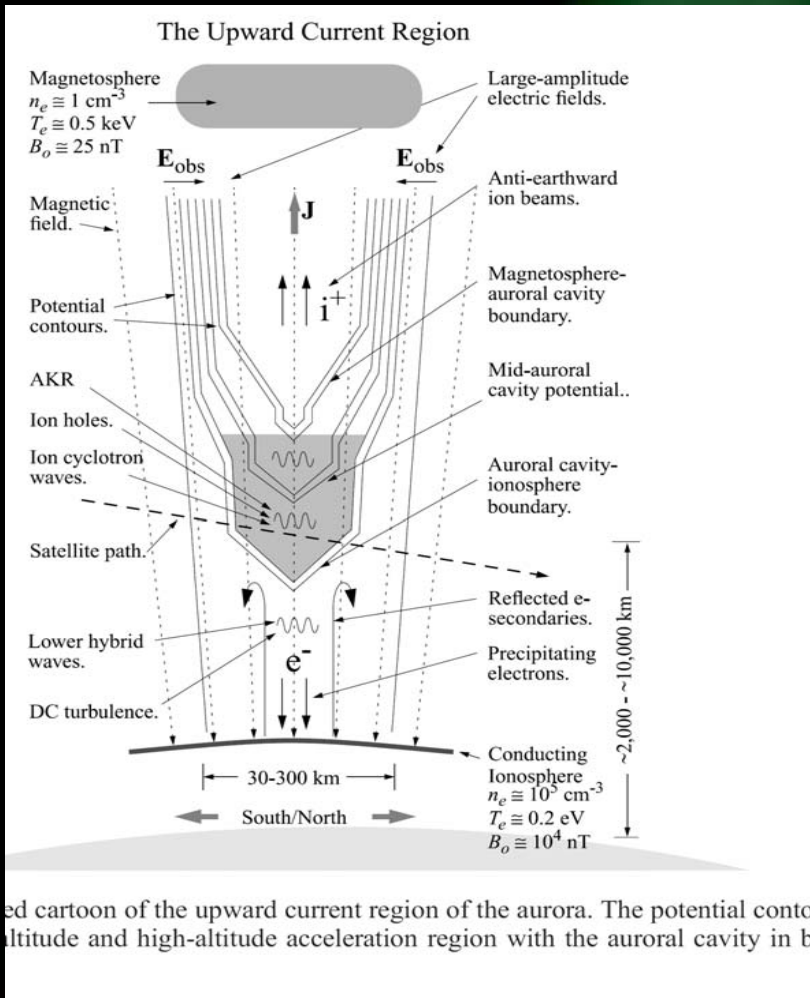


Model of cold beam producing secondaries (Evans, 1974)

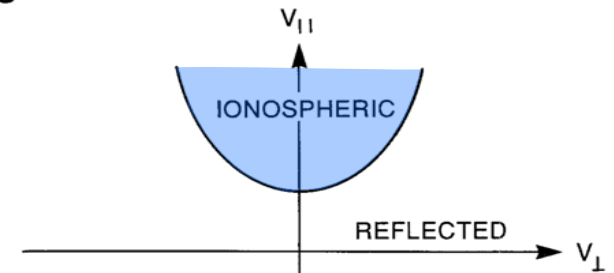


Model of hot electron beam and secondaries (Evans, 1974). Data from Franck and Ackerson, 1971

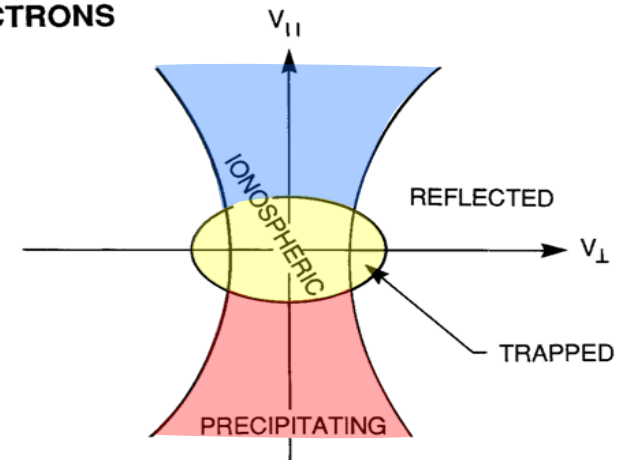
# Auroral cavity and trapped populations



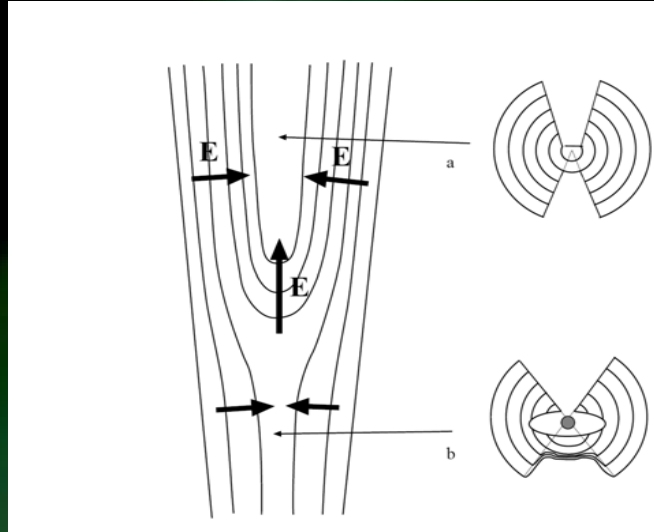
(a) IONS



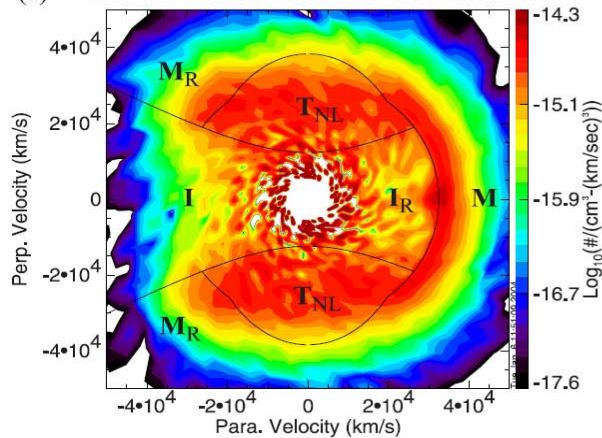
(b) ELECTRONS



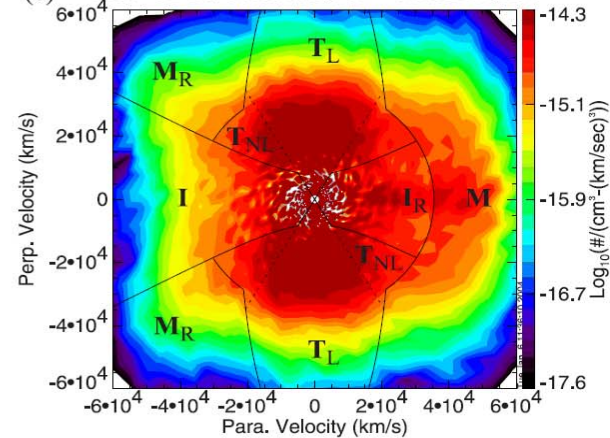
# Accelerated Maxwellian



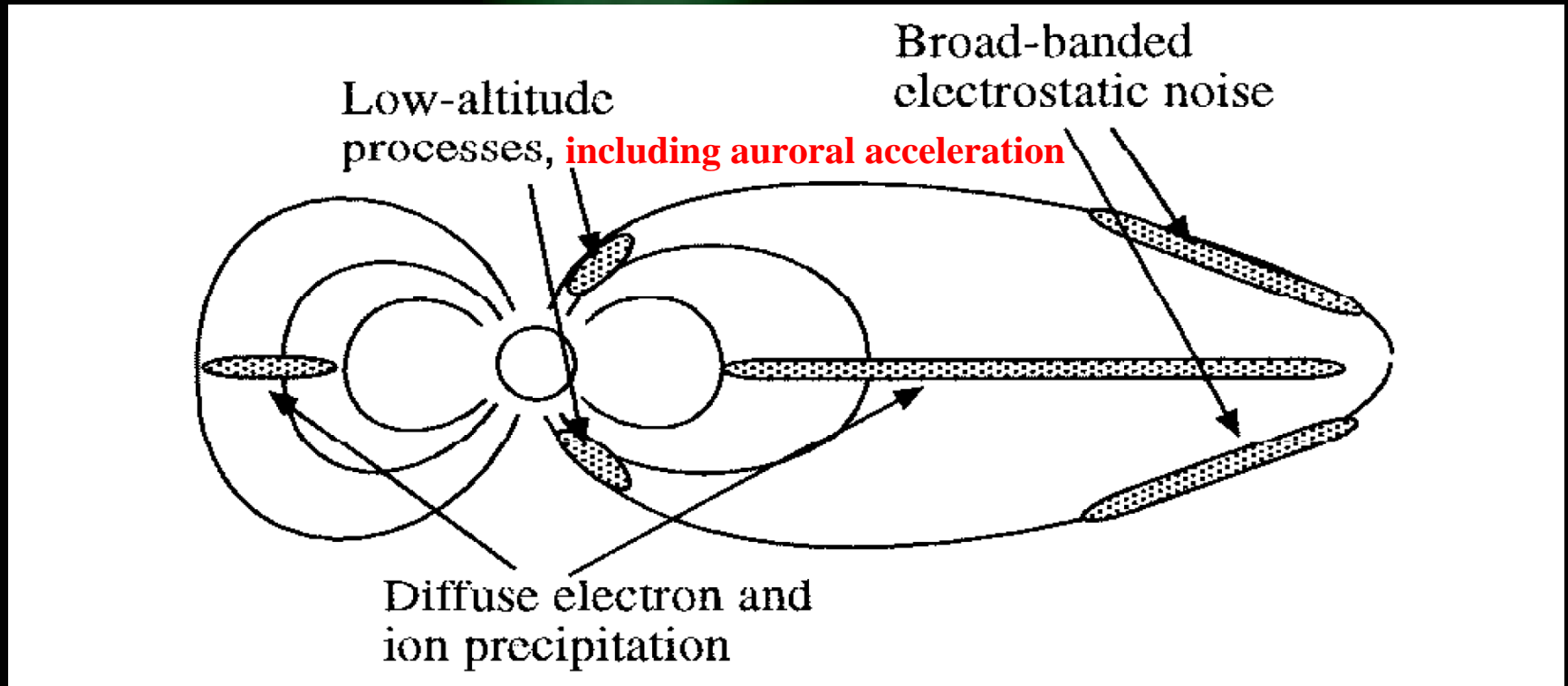
(a) FAST Eesa Survey DF  
1997-02-17/05:00:09.126 - 05:00:09.443



(c) FAST Eesa Survey DF  
1997-02-09/06:06:56.967 - 06:06:57.283



# Acceleration regions

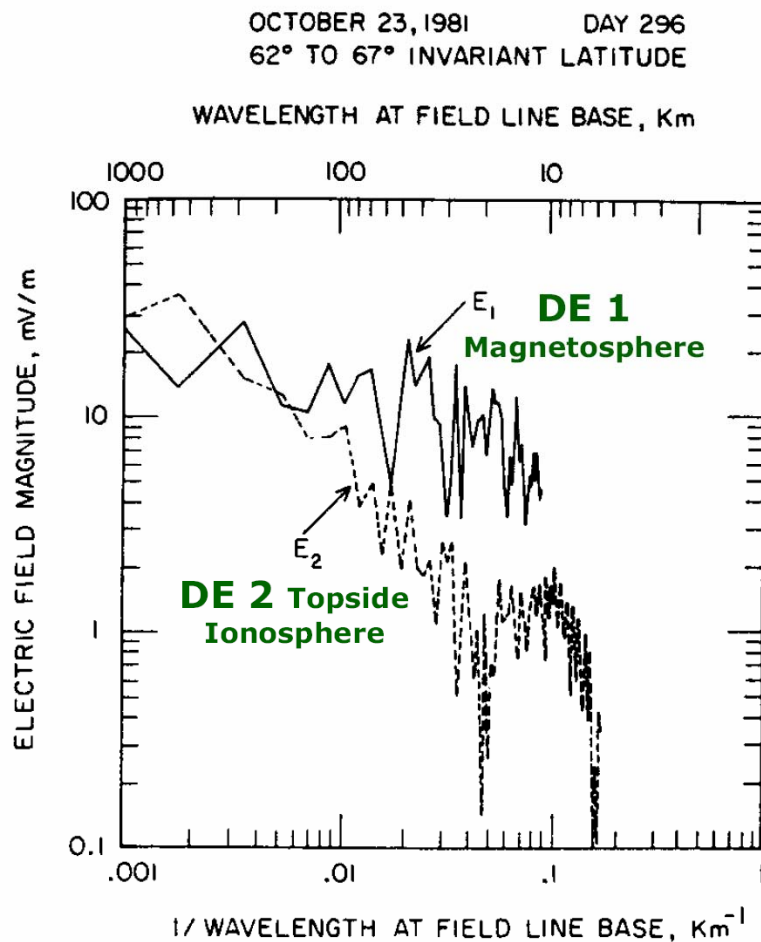


Koskinen

Auroral acceleration region typically situated at altitude of 1-3  $R_E$



# Mapping of auroral electric fields



Large  
scales

Small  
scales

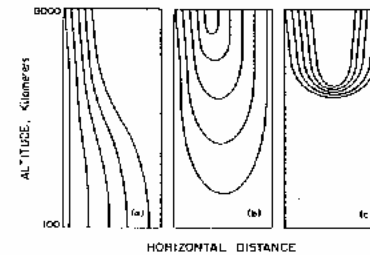


Figura 8. Contururi echipotentiale posibile in zona de accelerare (v. text). Din Hudson and Mozer (1978)

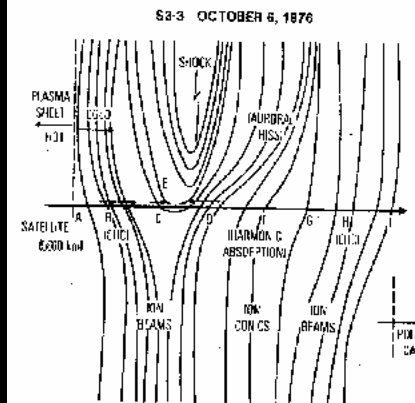


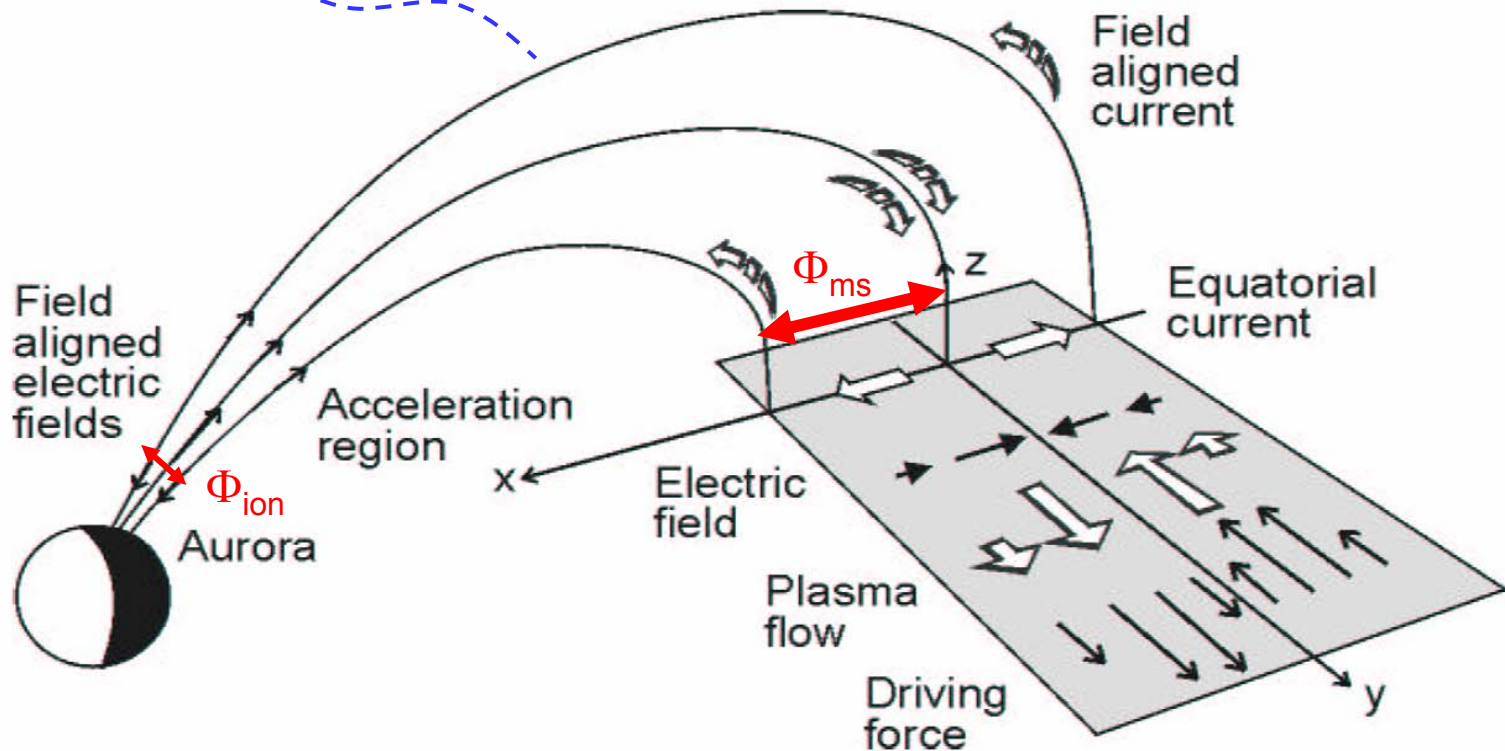
Figura 9. Reconstructia calitativa a contururilor echipotentiale din date S2-3. Figura ilustreaza 1 minut de observatii. Din Mizera et al. (1982)

Experimental results, comparison between Dynamics Explorer 1 and 2 at different altitudes. (Weimer et al., 1985)

# Static, medium-scale MI-coupling

## MI-coupling critical scale size II

Btw, mapping of E-field!



# Static, medium-scale MI-coupling

MI-coupling critical scale size III

$$j_{\parallel} = \Sigma_P \nabla_{\perp} \cdot \mathbf{E}_{\perp} + \mathbf{E}_{\perp} \cdot \nabla_{\perp} \Sigma_P + (\mathbf{E}_{\perp} \times \nabla_{\perp} \Sigma_H) \cdot \hat{\mathbf{b}}$$

$$j_{\parallel} = \Sigma_P \nabla_{\perp} \cdot \mathbf{E}_{\perp}$$

$$j_{\parallel} = K \Delta \Phi_{\parallel} = K(\Phi_{ms} - \Phi_{ion})$$

$$K(\Phi_{ms} - \Phi_{ion}) = \Sigma_P \nabla_{\perp} \cdot \mathbf{E}_{\perp ion}$$

$$K(\Phi_{ms} - \Phi_{ion}) = \Sigma_P \nabla_{\perp}^2 \Phi_{ion}$$

$$K(\Phi_{ms} - \Phi_{ion}) \approx \Sigma_P \frac{\Phi_{ion}}{L^2}$$

$$\Phi_{ion} = \left( 1 + \frac{\Sigma_P}{KL^2} \right)^{-1} \Phi_{ms}$$

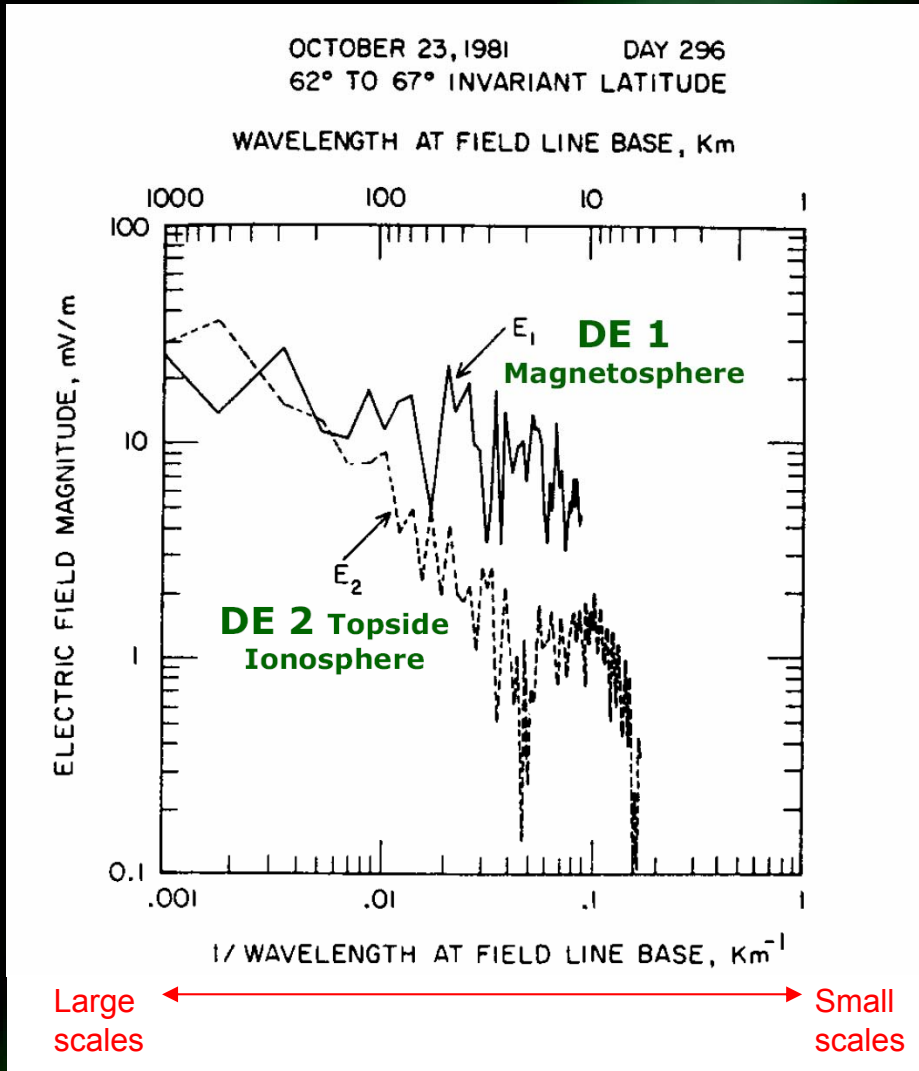
$$\text{When } L \gg \sqrt{\frac{\Sigma_P}{K}} : \quad \Phi_{ion} \approx \Phi_{ms}$$

$$\text{When } L \ll \sqrt{\frac{\Sigma_P}{K}} : \quad \Phi_{ion} \ll \Phi_{ms}$$

The last case means  $|\Phi_{ms}| \approx |\Delta \Phi_{\parallel}|$

# Static, medium-scale MI-coupling

## MI-coupling critical scale size



Experimental results, comparison between Dynamics Explorer 1 and 2 at different altitudes. (*Weimer et al., 1985*)

1/WAVELENGTH AT FIELD LINE BASE, Km<sup>-1</sup>

Fig. 4. Electric field spectrums from day 296 (October 23) of 1981. The spectrums are obtained from a Fourier transform of the electric field data between 62° and 67° invariant latitude. The solid line shows the spectrum of the electric field measured by DE 1. The solid line shows the spectrum of the electric field measured by DE 2. The ordinate values are obtained from the square root of the "spectral power density." The actual units are mV m<sup>-1</sup> km<sup>1/2</sup>. *Weimer et al, 1985*